



**B.Sc. (Sem. - 4) Physics**

**Course: US04CPHY21**

**Electromagnetic Theory and  
Spectroscopy**

**UNIT-3 Lecture 3**

# UNIT -III

## Atomic Spectra

# UNIT - III Atomic Spectra-Topics

- L-S Coupling
- J-J Coupling
- Fine structure of Hydrogen atom
- Spectral terms and their notations
- Positronium
- Mesonic atoms



# **Coupling Schemes:**

L-S Coupling

j-j Coupling

# Coupling Schemes: L-S Coupling

➔  $L = \sum l_i$  and  $S = \sum s_i$

➔  $J = L + S$

➔  $J = (l_1 + l_2 + l_3 + \dots) + (s_1 + s_2 + s_3 + \dots)$

➔ *J follows certain quantum rules*

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➔ *(i) All the three vectors L, S and J must be quantized.*

# Coupling Schemes: L-S Coupling

➤  $L = \sum l_i$  and  $S = \sum s_i$

➤  $J = L + S$

➤ *J follows certain quantum rules*

➤ *(ii) L may have values 0, 1, 2, 3, ... depending upon the **number of electrons** in the atom and **the directions of their orbital vectors.***



# Coupling Schemes: L-S Coupling

- ▶  $J$  follows certain quantum rules
- ▶ (iii)  $S$  may have values  $0, \frac{1}{2}, 1, \frac{3}{2}, 2, \frac{5}{2}, 3, \dots$  depending upon the **number of electrons** in the atom and the **directions of their spin vectors**.
- ▶  $S$  may have **integral** value for **even** number of electrons in the atom.

# Coupling Schemes: L-S Coupling

- *J follows certain quantum rules*
- *(iii) S may have values  $0, \frac{1}{2}, 1, \frac{3}{2}, 2, \frac{5}{2}, 3, \dots$  depending upon the number of electrons in the atom and the directions of their spin vectors.*
- *S may have **half-integral** for **odd** number of electrons in the atom.*

# Coupling Schemes: L-S Coupling

- ▶ *J follows certain quantum rules:*
- ▶ (iv) *J may have values 0, 1, 2, 3, ... when S is an integer*
- ▶ *i.e. for **even** number of electrons in the atom.*

# Coupling Schemes: L-S Coupling

- ▶ *J follows certain quantum rules*
- ▶ (iv) *J may have values  $\frac{1}{2}, \frac{3}{2}, \frac{5}{2}, \dots$  when  $S$  is a half-integer*
- ▶ *i.e. for **odd** number of electrons in the atom.*

# Coupling Schemes: L-S Coupling

- For ***an atom with only one electron*** in normal state
- $l = 0$
- $L = 0$  (Only one value)
- $S = \frac{1}{2}$
- $\therefore J = L + S = \frac{1}{2}$

# Coupling Schemes: L-S Coupling

➔ For an atom with **two electrons** in normal state

➔ Suppose  $l = 0$  and  $l = 1$

➔  $L = 0 + 1 = 1$  (Only one value)

➔  $S = \frac{1}{2} + \frac{1}{2} = 1$  or  $\frac{1}{2} - \frac{1}{2} = 0$

➔  $\therefore J = L + S = 0, 1, 2$  designated by  ${}^3P_0, {}^3P_1, {}^3P_2$

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Splitting of P state into a Triplet and  $J = 0$  and state  ${}^1P_1$

# Coupling Schemes: L-S Coupling

➔ Suppose  $L = 2$  and  $S = \frac{3}{2}$

Here  $L > S$  therefor there are  $(2S + 1)$  possible values of  $J$

➔  $\therefore J = L + S = (2 + \frac{3}{2} = \frac{7}{2}), \dots, (2 - \frac{3}{2} = \frac{1}{2})$

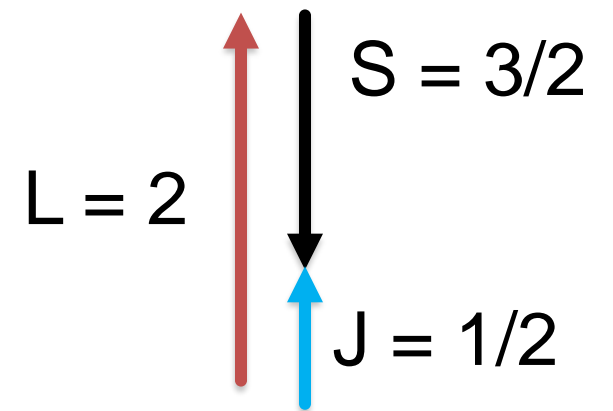
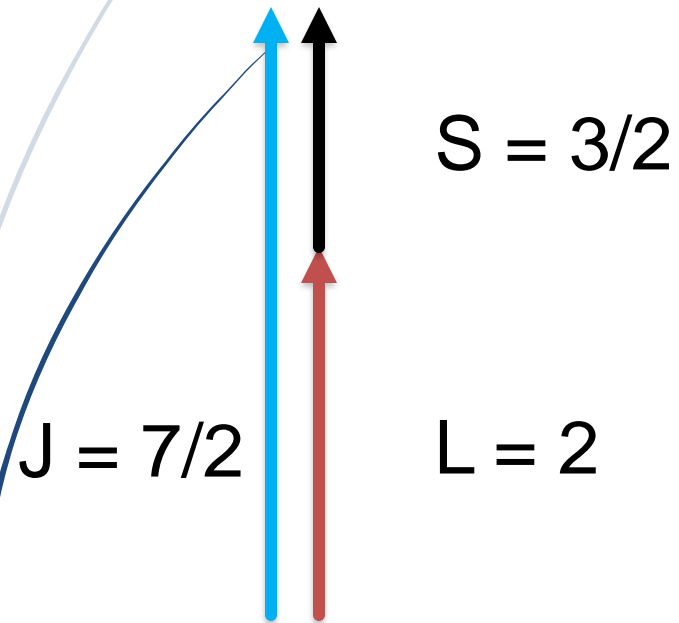
➔ Possible values of  $J$  are  $\frac{7}{2}, \frac{5}{2}, \frac{3}{2}, \frac{1}{2}$



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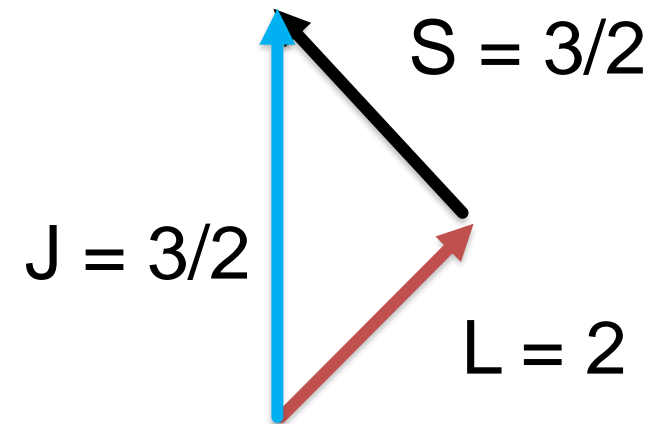
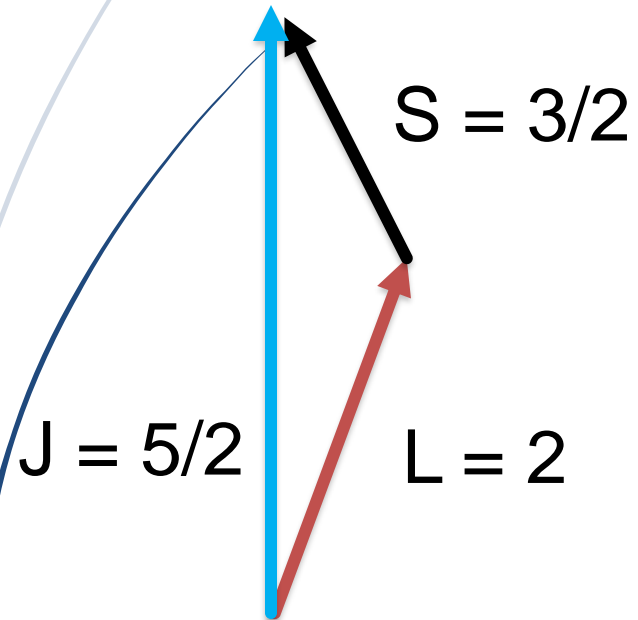
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# Coupling Schemes: L-S Coupling

➔ Suppose  $L = 2$  and  $S = \frac{3}{2}$

➔ Possible values of  $J$  are  $\frac{7}{2}, \frac{5}{2}, \frac{3}{2}, \frac{1}{2}$



# Coupling Schemes: L-S Coupling

➔ Suppose  $L = 1$  and  $S = 2$  [For practice]

Here  $L < S$  therefor there are  $(2L + 1)$  possible values of  $J$

➔ Possible values of  $J$  are ??????

# Coupling Schemes: j-j Coupling

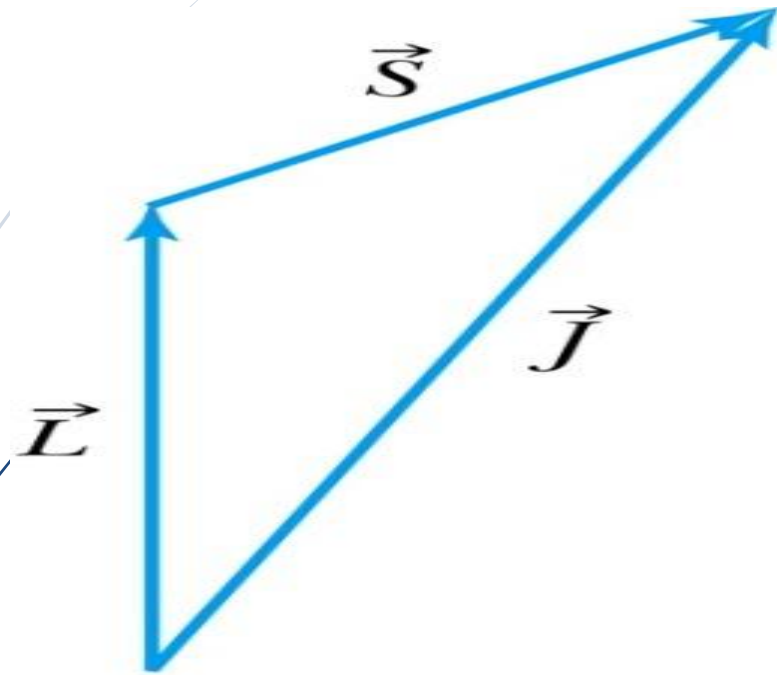
- ▶ In this type of coupling, each electron contributes to the total angular momentum of the atom by combining first its individual spin and orbit vectors by the relation

$$J = l + s$$

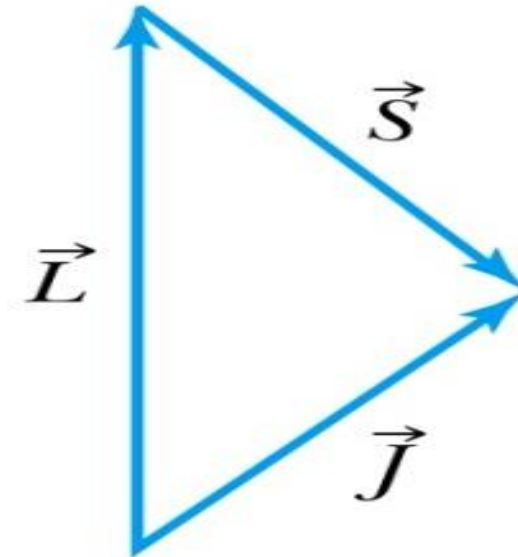
- ▶ Then the vector sum of each  $j$  corresponding to each electron gives the total angular momentum  $J$  to the atom

$$(l_1 + s_1) + (l_2 + s_2) + (l_3 + s_3) + \dots = J$$

# Coupling Schemes: j-j Coupling



$$j = \ell + s \\ = 1 + \frac{1}{2} = \frac{3}{2}$$



$$j = \ell - s \\ = 1 - \frac{1}{2} = \frac{1}{2}$$



## **3.7 FINE STRUCTURE OF HYDROGEN LINES**

## 3.7 FINE STRUCTURE OF HYDROGEN LINES

- How to explain fine structure doubling of spectral lines?
- Having explained the action of magnetic field on **small atomic magnets** in Stern-Gerlach experiment, we may visualize the **fine structure doubling of spectral lines** on the basis of a **magnetic interaction** between the **spin and angular momenta of atomic electrons**.

## 3.7 FINE STRUCTURE OF HYDROGEN LINES

- **Magnetic Moment:**
- An **electron** revolving round a **proton** finds itself in a magnetic field because, in its own frame of reference, the proton is circling about it.
- This magnetic field then acts upon the electron's own **magnetic moment** to produce a kind of internal Zeeman effect.



## 3.7 FINE STRUCTURE OF HYDROGEN LINES

- The magnetic energy  $V_m$  of a magnetic dipole moment  $\mu$  in a magnetic field of flux density  $\mathbf{B}$  is

$$V_m = -\mu \cdot \mathbf{B} \quad (1)$$

$$V_m = \pm \frac{e \hbar}{2 m} \mathbf{B} \quad (2)$$

- where  $\frac{e \hbar}{2 m}$  is **spin magnetic moment** in the direction of  $\mathbf{B}$ .

## 3.7 FINE STRUCTURE OF HYDROGEN LINES

- Depending upon the **orientation of its spin vector**, the energy of the electron in a given atomic state will be **higher** or **lower** by the  $\frac{e \hbar}{2 m} \mathbf{B}$  than its value in the absence of spin orbit interaction.

## 3.7 FINE STRUCTURE OF HYDROGEN LINES

- The result is the splitting of every quantum state (except S state) into two separate sub-states and consequently, the **splitting of every spectral line into two component lines.**

## 3.7 FINE STRUCTURE OF HYDROGEN LINES

- Depending upon the **orientation of its spin vector**, the energy of the electron in a given atomic state will be **higher or lower** by the  $\frac{e \hbar}{2 m} B$  than its value in the absence of spin orbit interaction.
- The result is the splitting of every quantum state (except S state) into two separate sub-states and, consequently, the **splitting of every spectral line into two component lines.**
- The assignment of  $s = \frac{1}{2}$  is the **only one** that conforms to the **observed fine structure doubling.**

## 3.7 FINE STRUCTURE OF HYDROGEN LINES

- Let us now have the estimate of **magnetic flux density** and **magnetic energy** in fine structure doubling.
- The circular wire loop of radius  $r$  that carries the current  $i$  has a magnetic field of flux density at its centre

$$B = \frac{\mu_0 i}{2 r} \quad (3)$$

- where  $\mu_0$  is the permeability of the vacuum.

## 3.7 FINE STRUCTURE OF HYDROGEN LINES

- An **orbital electron** say in a **hydrogen atoms** 'sees' itself circled  **$f$**  times each second by a proton of charge  **$+e$**  for a resulting **flux density** of

$$B = \frac{\mu_0 f e}{2 r} \quad (4)$$

- In the **ground state of hydrogen atom**  **$f = 6.8 \times 10^{15}$  cps** and  **$r = 5.3 \times 10^{-11}$  metre**, so that

$$B = 13 \text{ Tesla (abbreviation T)} \quad (5)$$

- which is a **very strong magnetic field**.

## 3.7 FINE STRUCTURE OF HYDROGEN LINES

- The value of **Bohr magneton** is

$$\frac{e \hbar}{2 m} = 9.27 \times 10^{-24} \text{ J/T} \quad (5)$$

- Hence, **magnetic energy** (  $V_m$  ) of one such electron is

$$V_m = \frac{e \hbar}{2 m} B = ( 9.27 \times 10^{-24} \text{ J/T} ) \times 13 \text{ T}$$

$$V_m = 1.2 \times 10^{-22} \text{ joules} \quad (6)$$

## 3.7 FINE STRUCTURE OF HYDROGEN LINES

- The **wavelength shift** in such a change in energy is about **2  $^{\circ}\text{A}$**  for a spectral line of unperturbed wavelength **6563  $^{\circ}\text{A}$** , somewhat **more than the observed splitting of the line** originating in the  $n = 3 \rightarrow n = 2$  transition.
- However, the **flux density of the magnetic field at the orbits of higher order is less than for ground state orbit**, which accounts for the discrepancy.





## **3.8. SPECTRAL TERMS AND THEIR NOTATIONS**

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- The optical spectrum of an element is the characteristic of the valence electron, i.e., optical spectrum particularly depends upon the electrons which are not interlocked in closed shells.

## 3.8. SPECTRAL TERMS AND THEIR NOTATIONS

- The different atoms having different number of valence electrons in their outermost orbit show different types of spectra.

## 3.8. SPECTRAL TERMS AND THEIR NOTATIONS

- For example, the **alkali metals** show one type of spectrum, while **alkaline earth of another type**, since they possess **one** and **two** valence electrons, respectively in their outermost shell.

## 3.8. SPECTRAL TERMS AND THEIR NOTATIONS

- Further, the **state of electron**, is described in terms of the different values of  $l$ ,  $s$  and  $j$ .
  - The **small letters**  $l$ ,  $s$  and  $j$  depict the **state of an electron**, while
  - the **capital letters**,  $L$ ,  $S$  and  $J$  depict the **state of an atom as a whole**

## 3.8. SPECTRAL TERMS AND THEIR NOTATIONS

- ▶ But when we are dealing with **one valence electron system**, the values of  $L$ ,  $S$  and  $J$  are the **same** as those of  $l$ ,  $s$  and  $j$ , respectively, since the **interlocked electrons in closed shells and sub-shells contribute nothing to the total angular momentum.**

## 3.8. SPECTRAL TERMS AND THEIR NOTATIONS

- Dealing with **many electron systems**, the vectors  $L$ ,  $S$  and  $J$  which define the **state of an atom** are vector sums of  $l$ ,  $s$  and  $j$ , respectively for different free electrons.

## 3.8. SPECTRAL TERMS AND THEIR NOTATIONS

- The electrons having 0, 1, 2, 3.. etc. values for  $l$  are represented by the small letters s, p, d, f, ..... etc. respectively.
- Similarly, the capital letters, S, P, D, F depict the state of the atom for the value of  $L$  as 0, 1, 2, 3, ...,etc.



## 3.8. SPECTRAL TERMS AND THEIR NOTATIONS

- In the case of **one-electron system** the value of  $S$  is  $+\frac{1}{2}$  and hence the **multiplicity of the state** is 2.
- The multiplicity of a state is given by  $r = 2S + 1$ .
- Hence the **single-electron system always gives rise to a double state**, corresponding to the value  $(L + \frac{1}{2})$  and  $(L - \frac{1}{2})$  for  $J$  with the exception of ground state.

$$J = (L + \frac{1}{2}) \text{ and } J = (L - \frac{1}{2})$$

## 3.8. SPECTRAL TERMS AND THEIR NOTATIONS

- The multiplicity of a state is given by  $r = 2S + 1$ .
- In the case of **many electron system**,  $S$  may not necessarily have the value  $1/2$ , but, it may have any value.
- For example, in **two electron system**  $S = 0$  or  $1$ . Therefore, the state is either **singlet** or **triplet**.

## 3.8. SPECTRAL TERMS AND THEIR NOTATIONS

- The multiplicity of a state is given by  $r = 2S + 1$ .
- In **three electron system**  $S = \frac{1}{2}$  or  $\frac{3}{2}$  and hence state is **doublet** or **quartet**, and so on except the ground term.

## 3.8. SPECTRAL TERMS AND THEIR NOTATIONS

- Now if we want to depict the **state of an atom**, we must mention the values of L, S, J and r.
- This is done by writing the state having L=1, S =1/2 and as J = 3/2 as  ${}^2P_{3/2}$ .
- The term  ${}^2P_{3/2}$ , clearly indicates that
  - the value of L is given by the capital letter,
  - the value of J by right subscript below the capital letter and
  - the value of r by left superscript at the top of capital letter.
- Since  $r = 2S + 1$ , the S is obviously determined.

## 3.8. SPECTRAL TERMS AND THEIR NOTATIONS

- $L=1$ ,  $S = 1/2$  and as  $J = 3/2$  as  ${}^2P_{3/2}$ .
- Sometimes the value of principal quantum number is also mentioned before the capital letter as  $2 {}^2P_{3/2}$
- Here the value of  $n$  is 2.

## 3.8. SPECTRAL TERMS AND THEIR NOTATIONS

- In **one-electron system**, every state is **doublet** and may be depicted as below:

$${}^2S_{1/2}, {}^2P^{1/2}, {}^2P_{3/2}, {}^2D_{3/2}, {}^2D_{5/2} \text{ and so on.}$$

- Here one point is important to note that **multiplicity symbol of the system is used always** whether all the terms are present or not.
- For example, the ground term is here written as  ${}^2S_{1/2}$ . But actually, the state is singlet and we should write  ${}^1S_{1/2}$ .

## 3.8. SPECTRAL TERMS AND THEIR NOTATIONS

- To write the **ground state** as  ${}^2S_{1/2}$  is preferable because it indicates to which system the ground term belongs, as in this case to the doublet system.

## 3.8. SPECTRAL TERMS AND THEIR NOTATIONS

- **Why is the ground state always singlet?**
- We know that the value of  $J$  is given by  $L + S$  to  $L - S$  with a difference of one.
- Since in **one-electron system**  $S = \pm 1/2$  (with respect to  $L$ ) hence  $J = L + 1/2$  and  $J = L - 1/2$ , i.e.  $J$  has only two values and the state is said to be **doublet**.



## 3.8. SPECTRAL TERMS AND THEIR NOTATIONS

- For the ground state  $L = 0$ ,  $J = \frac{1}{2}$  or  $J = -\frac{1}{2}$ . But **net angular momentum of the atom is always positive** and the possibility of  $-1/2$  is ruled out.
- Therefore, the ground state of single electron system is always singlet.

## 3.8. SPECTRAL TERMS AND THEIR NOTATIONS

- In case of **many-electron system** for the ground state  $L = 0$  and  $J = S$ ,  $S$  can assume any value of  $0, 1/2, 3/2$  etc.
- When  $L < S$ , the multiplicity of the state is given by  $(2L + 1)$  which leads again to the possible value of  $J$  as **one** and the state is **singlet**.

## 3.8. SPECTRAL TERMS AND THEIR NOTATIONS

- The notation for a single-electron atom becomes  $n^{2S+1} L_J$
- The letters and numbers are called **spectroscopic symbols**.
- There are **singlet** states ( $S = 0$ ) and **triplet** states ( $S = 1$ ) for two electrons.



## **3.9 POSITRONIUM**

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- The combined structure, the **positronium**, as it is called, revolves round the common centre of gravity.
- Thus the system is very **much identical to the hydrogen atom** and so the **formulae derived for hydrogen atom may be applied for this system.**

## 3.9 POSITRONIUM

- Hence the **orbital radius** is given by

$$r_n = \frac{\hbar^2}{e^2} \frac{m+m}{m^2} \frac{n^2}{Z} (4 \pi \epsilon_0) \quad (1)$$

- For ground state,  $r_0 = \frac{2 \hbar^2 (4 \pi \epsilon_0)}{m e^2} \approx 1 \text{ }^\circ\text{A}$
- Thus, the **orbital radius increases by a factor two.**

## 3.9 POSITRONIUM

- The **energy of dissociation** is given by

$$E = - \frac{1}{(4 \pi \epsilon_0)^2} \frac{e^4}{2 \hbar^2} \frac{m^2}{m+m} = - \frac{m e^4}{4 \hbar^2} \frac{1}{(4 \pi \epsilon_0)^2} \quad (2)$$

- This is **one half of the ionization potential of the hydrogen atom.**

## 3.9 POSITRONIUM

- The positronium exists in **two states** depending upon the state of spin orientation.
- These are **ortho-positronium** in which the spins of the two particles are **parallel** and **para-positronium** in which spins are **anti-parallel**.



## 3.9 POSITRONIUM

- In **para-positronium**, as two particles have their spin **anti-parallel**, the system is an **unstable** one with a life time  $1.25 \times 10^{-10}$  sec. The para-positronium decays into **two** photons.
- The **ortho-positronium** is **stable** enough and has a life time of the order of  $1.4 \times 10^{-7}$  sec.
- It decays into **three**  $\gamma$ -photons so as to maintain the conservation of spin.

## 3.9 POSITRONIUM

- The ground level of **ortho-positronium** lies above the ground state of para-positronium by only  $0.84 \times 10^{-2}$  eV.
- The difference between the life times of two photons accelerates the annihilation process.



➔ Thank you..

